

Objectives

1. Derivation of mass, momentum and energy conservation using Boltzmann's equation.

Rewriting again for the Boltzmann's Equation, we get:

$$\frac{\partial f_j}{\partial t} + \mathbf{v}_j \cdot \nabla_{\mathbf{r}} f_j + \frac{\mathbf{F}_j}{m_j} \cdot \nabla_{\mathbf{v}_j} f_j = 2\pi \sum_i \int_{-\infty}^{\infty} \int_0^{\infty} (f'_i f'_j - f_i f_j) v_r b db d\mathbf{v}_i \quad (1)$$

Consider a property of a particle ϕ_j . ϕ_j can be either m_j , $m_j \mathbf{v}_j$ or $\frac{mv_j^2}{2}$. Let us first multiply both sides of this equation by ϕ_j and then integrate with respect to $\mathbf{v}_j$ and then sum over all j . For right hand side, we have:

$$2\pi \sum_j \sum_i \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_0^{\infty} \phi_j (f'_i f'_j - f_i f_j) v_r b db d\mathbf{v}_i d\mathbf{v}_j$$

By inspection, we can know that:

$$\begin{aligned}
& \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_0^{\infty} \phi_j (f'_i f'_j - f_i f_j) v_r b d b d \mathbf{v}_i d \mathbf{v}_j \\
&= - \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_0^{\infty} \phi'_j (f'_i f'_j - f_i f_j) v_r b d b d \mathbf{v}_i d \mathbf{v}_j \\
&= \frac{1}{2} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_0^{\infty} (\phi_j - \phi'_j) (f'_i f'_j - f_i f_j) v_r b d b d \mathbf{v}_i d \mathbf{v}_j \\
& \sum_i \sum_j \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_0^{\infty} (\phi_j - \phi'_j) (f'_i f'_j - f_i f_j) v_r b d b d \mathbf{v}_i d \mathbf{v}_j
\end{aligned}$$

The double sum operators allow us to exchange the indices i and j , and thus allowing us to write:

$$\begin{aligned}
& \sum_i \sum_j \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_0^{\infty} (\phi_j - \phi'_j) (f'_i f'_j - f_i f_j) v_r b d b d \mathbf{v}_i d \mathbf{v}_j \\
&= \sum_i \sum_j \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_0^{\infty} (\phi_i - \phi'_i) (f'_i f'_j - f_i f_j) v_r b d b d \mathbf{v}_i d \mathbf{v}_j
\end{aligned}$$

$$\begin{aligned}
& \sum_i \sum_j \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_0^{\infty} \phi_j(f'_i f'_j - f_i f_j) v_r b db d\mathbf{v}_i d\mathbf{v}_j \\
&= \frac{1}{4} \sum_i \sum_j \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_0^{\infty} (\phi_j + \phi_i - \phi'_j - \phi'_i)(f'_i f'_j - f_i f_j) v_r b db d\mathbf{v}_i d\mathbf{v}_j \\
&= 0
\end{aligned}$$

This is because $\phi_j + \phi_i - \phi'_j - \phi'_i = 0$. Before any more derivation, bear in mind that:

$$\mathbf{v}_0 = \frac{\sum_j \rho_j m_j \langle \mathbf{v}_j \rangle}{\sum_j \rho_j m_j} \quad (2)$$

which means:

$$\rho_m \mathbf{v}_0 = \sum_j \rho_j m_j \langle \mathbf{v}_j \rangle$$

$$\rho_m = \sum_j \rho_j m_j.$$

For the left hand side, if we neglect the force term, we have:

$$\sum_j \int_{-\infty}^{\infty} \phi_j \frac{\partial f_j}{\partial t} d\mathbf{v}_j + \sum_j \int_{-\infty}^{\infty} \phi_j \mathbf{v}_j \cdot \nabla_{\mathbf{r}} f_j d\mathbf{v}_j$$

Bear in mind that ϕ_j is a function of the coordinate $\mathbf{v}_j$, which is independent of $\mathbf{r}$ and t . Therefore, we can write:

$$\sum_j \int_{-\infty}^{\infty} \frac{\partial}{\partial t} (\phi_j f_j) d\mathbf{v}_j + \sum_j \int_{-\infty}^{\infty} \nabla_{\mathbf{r}} \cdot (\mathbf{v}_j f_j \phi_j) d\mathbf{v}_j$$

Putting all these together, we have:

$$\sum_j \int_{-\infty}^{\infty} \frac{\partial}{\partial t} (\phi_j f_j) d\mathbf{v}_j + \sum_j \int_{-\infty}^{\infty} \nabla_{\mathbf{r}} \cdot (\mathbf{v}_j f_j \phi_j) d\mathbf{v}_j = 0$$

The easiest case is that when $\phi_j = m_j$.

$$\frac{\partial}{\partial t} \sum_j m_j \int_{-\infty}^{\infty} f_j d\mathbf{v}_j + \nabla_{\mathbf{r}} \cdot \sum_j \int_{-\infty}^{\infty} m_j (\mathbf{v}_j f_j) d\mathbf{v}_j = 0$$

$$\frac{\partial}{\partial t} \sum_j \rho_j m_j + \nabla_{\mathbf{r}} \cdot \sum_j \rho_j m_j \langle \mathbf{v}_j \rangle = 0$$

$$\frac{\partial \rho_m}{\partial t} + \nabla_{\mathbf{r}} \cdot (\rho_m \mathbf{v}_0) = 0$$

If we have $\phi_j = m_j \mathbf{v}_j$. Let us neglect the summation operator.

$$\int_{-\infty}^{\infty} m_j \mathbf{v}_j \frac{\partial f_j}{\partial t} d\mathbf{v}_j + \int_{-\infty}^{\infty} m_j \mathbf{v}_j \mathbf{v}_j \cdot \nabla_{\mathbf{r}} f_j d\mathbf{v}_j = 0$$

$$\frac{\partial}{\partial t} m_j \int_{-\infty}^{\infty} \mathbf{v}_j f_j d\mathbf{v}_j + \nabla_{\mathbf{r}} \cdot m_j \int_{-\infty}^{\infty} \mathbf{v}_j \mathbf{v}_j f_j d\mathbf{v}_j = 0$$

Attention must be drawn to the second term:

$$\mathbf{v}_j \mathbf{v}_j = (\mathbf{v}_j - \mathbf{v}_0)(\mathbf{v}_j - \mathbf{v}_0) + 2\mathbf{v}_j \mathbf{v}_0 - \mathbf{v}_0 \mathbf{v}_0$$

The pressure tensor is defined as:

$$\mathbf{P}_j = m_j \int_{-\infty}^{\infty} (\mathbf{v}_j - \mathbf{v}_0)(\mathbf{v}_j - \mathbf{v}_0) f_j d\mathbf{v}_j$$

Then we get:

$$\frac{\partial}{\partial t}(\rho_j m_j \langle \mathbf{v}_j \rangle) + \nabla_{\mathbf{r}} \cdot \mathbf{P}_j + \nabla_{\mathbf{r}} \cdot 2\rho_j m_j \langle \mathbf{v}_j \rangle \mathbf{v}_0 - \nabla_{\mathbf{r}} \cdot (\rho_j m_j \mathbf{v}_0 \mathbf{v}_0) = 0$$

Sum over all j .

$$\frac{\partial}{\partial t}(\rho_m \mathbf{v}_0) + \nabla_{\mathbf{r}} \cdot \mathbf{P} + \nabla_{\mathbf{r}} \cdot 2\rho_m \mathbf{v}_0 \mathbf{v}_0 - \nabla_{\mathbf{r}} \cdot (\rho_m \mathbf{v}_0 \mathbf{v}_0) = 0$$

$$\begin{aligned}\nabla_{\mathbf{r}} \cdot (\rho_m \mathbf{v}_0 \mathbf{v}_0) &= \frac{\partial}{\partial k} \rho_m v_{0,k} v_{0,l} = \rho_m v_{0,k} \frac{\partial}{\partial k} v_{0,l} + v_{0,l} \frac{\partial}{\partial k} \rho_m v_{0,k} \\ &= \rho_m \mathbf{v}_0 \cdot \nabla_{\mathbf{r}} \mathbf{v}_0 + \mathbf{v}_0 \nabla_{\mathbf{r}} \cdot (\rho_m \mathbf{v}_0)\end{aligned}$$

$$\begin{aligned}\mathbf{v}_0 \left[\frac{\partial \rho_m}{\partial t} + \nabla_{\mathbf{r}} \cdot (\rho_m \mathbf{v}_0) \right] + \rho_m \left[\frac{\partial \mathbf{v}_0}{\partial t} + \mathbf{v}_0 \cdot \nabla_{\mathbf{r}} \mathbf{v}_0 \right] + \nabla_{\mathbf{r}} \cdot \mathbf{P} &= 0 \\ \rho_m \left[\frac{\partial \mathbf{v}_0}{\partial t} + \mathbf{v}_0 \cdot \nabla_{\mathbf{r}} \mathbf{v}_0 \right] + \nabla_{\mathbf{r}} \cdot \mathbf{P} &= 0\end{aligned}$$

Finally, to derive the energy balance equation, the property $\frac{1}{2}m_j v_j^2$ has to be considered. Without any integration and summation operator,

$$\frac{1}{2}m_j v_j^2 \frac{\partial f_j}{\partial t} + \frac{1}{2}m_j v_j^2 \mathbf{v}_j \cdot \nabla_{\mathbf{r}} f_j = 0$$

which can be rewritten as:

$$\frac{\partial}{\partial t} \left[\frac{1}{2} f_j m_j (v_j - v_0)^2 + \frac{1}{2} f_j m_j v_0^2 \right] + \nabla_{\mathbf{r}} \cdot f_j \mathbf{v}_j \frac{1}{2} m_j v_j^2 = 0 \quad (3)$$

Looking at the second term:

$$v_j^2 \mathbf{v}_j = (v_j - v_0)^2 (\mathbf{v}_j - \mathbf{v}_0) + (v_j - v_0)^2 \mathbf{v}_0 + 2(\mathbf{v}_j \cdot \mathbf{v}_0) \mathbf{v}_j - v_0^2 \mathbf{v}_j \quad (4)$$

Tackle them one by one: We know that:

$$\mathbf{q}_j = \frac{m_j}{2} \int_{-\infty}^{\infty} (v_j - v_0)^2 (\mathbf{v}_j - \mathbf{v}_0) f_j d\mathbf{v}_j$$

$$\rho_j \mathbf{v}_0 \tilde{E}_j = \frac{\mathbf{v}_0}{2} \int_{-\infty}^{\infty} (v_j - v_0)^2 f_j d\mathbf{v}_j$$

Then, after applying integration over all $\mathbf{v}_j$:

$$\begin{aligned} \frac{\partial}{\partial t} \left[\rho_j m_j \tilde{E}_j + \frac{\rho_j m_j v_0^2}{2} \right] + \nabla_{\mathbf{r}} \cdot \left[\mathbf{q}_j + \rho_j m_j \tilde{E}_j \mathbf{v}_0 - \frac{1}{2} \rho_j m_j \langle \mathbf{v}_j \rangle v_0^2 \right. \\ \left. + m_j \int f_j (\mathbf{v}_j \cdot \mathbf{v}_0) \mathbf{v}_j d\mathbf{v}_j \right] = 0 \end{aligned}$$

That integral has to be taken care of. Thinking in terms of index notation:

$$\begin{aligned} \sum_{k=x,y,z} v_{0,k} (v_{j,k} - v_{0,k}) (v_{j,l} - v_{0,l}) \\ = v_{0,k} (v_{j,k} v_{j,l} - v_{j,k} v_{0,l} - v_{0,k} v_{j,l} + v_{0,k} v_{0,l}) \end{aligned}$$

This gives us that:

$$\mathbf{v}_0 \cdot (\mathbf{v}_j - \mathbf{v}_0)(\mathbf{v}_j - \mathbf{v}_0) = (\mathbf{v}_j \cdot \mathbf{v}_0)\mathbf{v}_j - (\mathbf{v}_j \cdot \mathbf{v}_0)\mathbf{v}_0 - v_0^2\mathbf{v}_j + v_0^2\mathbf{v}_0$$

We obtain the relation:

$$(\mathbf{v}_j \cdot \mathbf{v}_0)\mathbf{v}_j = (\mathbf{v}_j \cdot \mathbf{v}_0)\mathbf{v}_0 + \mathbf{v}_0 \cdot (\mathbf{v}_j - \mathbf{v}_0)(\mathbf{v}_j - \mathbf{v}_0) + v_0^2\mathbf{v}_j - v_0^2\mathbf{v}_0 \quad (5)$$

$$m_j \left[\int f_j \mathbf{v}_j d\mathbf{v}_j \right] \cdot \mathbf{v}_0 \mathbf{v}_0 = \rho_j m_j \langle \mathbf{v}_j \rangle \cdot \mathbf{v}_0 \mathbf{v}_0$$

$$\mathbf{v}_0 \cdot m_j \int f_j (\mathbf{v}_j - \mathbf{v}_0)(\mathbf{v}_j - \mathbf{v}_0) d\mathbf{v}_j = \mathbf{v}_0 \cdot \mathbf{P}_j$$

$$m_j v_0^2 \int f_j \mathbf{v}_j d\mathbf{v}_j = \rho_j m_j \langle \mathbf{v}_j \rangle v_0^2$$

$$m_j v_0^2 \mathbf{v}_0 \int f_j d\mathbf{v}_j = \rho_j m_j v_0^2 \mathbf{v}_0$$

Now sum over all these terms over j , we obtain:

$$\begin{aligned} \sum_j \left[\rho_j m_j \langle \mathbf{v}_j \rangle \cdot \mathbf{v}_0 \mathbf{v}_0 + \mathbf{v}_0 \cdot \mathbf{P}_j + \rho_j m_j \langle \mathbf{v}_j \rangle v_0^2 - \rho_j m_j v_0^2 \mathbf{v}_0 \right] \\ = \rho_m v_0^2 \mathbf{v}_0 + \mathbf{v}_0 \cdot \mathbf{P} \end{aligned}$$

The entire equation becomes:

$$\begin{aligned} \frac{\partial \rho_m \tilde{E}}{\partial t} + \frac{1}{2} \frac{\partial \rho_m v_0^2}{\partial t} + \nabla_{\mathbf{r}} \cdot \left[\mathbf{q} + \rho_m \tilde{E} \mathbf{v}_0 - \frac{\rho_m v_0^2}{2} \mathbf{v}_0 + \rho_m v_0^2 \mathbf{v}_0 + \mathbf{v}_0 \cdot \mathbf{P} \right] \\ = 0 \end{aligned}$$

Look for mass balance relation and material derivative:

$$\begin{aligned} \tilde{E} \frac{\partial \rho_m}{\partial t} + \tilde{E} \nabla_{\mathbf{r}} \cdot (\rho_m \mathbf{v}_0) &= 0 \\ \rho_m \frac{\partial \tilde{E}}{\partial t} + \rho_m \mathbf{v}_0 \cdot \nabla_{\mathbf{r}} \tilde{E} &= \rho_m \frac{D\tilde{E}}{Dt} \\ \frac{1}{2} \left[v_0^2 \frac{\partial \rho_m}{\partial t} + v_0^2 \nabla_{\mathbf{r}} \cdot \rho_m \mathbf{v}_0 \right] &= 0 \end{aligned}$$

$$\rho_m \frac{D\tilde{E}}{Dt} + \frac{\rho_m}{2} \frac{\partial v_0^2}{\partial t} + \nabla_{\mathbf{r}} \cdot \mathbf{q} + \frac{\rho_m \mathbf{v}_0}{2} \cdot \nabla_{\mathbf{r}} v_0^2 + \nabla_{\mathbf{r}} \cdot \mathbf{v}_0 \cdot \mathbf{P} = 0$$

$$\rho_m \frac{D\tilde{E}}{Dt} + \rho_m \mathbf{v}_0 \cdot \frac{\partial \mathbf{v}_0}{\partial t} + \nabla_{\mathbf{r}} \cdot \mathbf{q} + \rho_m \mathbf{v}_0 \mathbf{v}_0 : \nabla_{\mathbf{r}} \mathbf{v}_0 + \nabla_{\mathbf{r}} \cdot \mathbf{v}_0 \cdot \mathbf{P} = 0$$

Multiply both sides of the momentum balance by $\mathbf{v}_0$.

$$\rho_m \mathbf{v}_0 \cdot \frac{\partial \mathbf{v}_0}{\partial t} + \rho_m \mathbf{v}_0 \mathbf{v}_0 : \nabla_{\mathbf{r}} \mathbf{v}_0 + \mathbf{v}_0 \cdot \nabla_{\mathbf{r}} \cdot \mathbf{P} = 0$$

Then, we get:

$$\rho_m \frac{D\tilde{E}}{Dt} + \nabla_{\mathbf{r}} \cdot \mathbf{q} + \mathbf{P} : \nabla_{\mathbf{r}} \mathbf{v}_0 = 0$$